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Tightness of Jensen’s Bounds and Applications to MIMO Communications

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Abstract—Due to the difficulty in manipulating the distribution of Wishart random matrices, the performance analysis of multiple-input-multiple-output (MIMO) channels has mainly focused on deriving capacity bounds via Jensen’s inequality. However, to the best of our knowledge, the tightness of Jensen’s bounds has not yet been rigorously quantified in the general MIMO context. This paper proposes a new methodology for measuring the tightness of Jensen’s bounds via the sandwich theorem. In particular, we first compare the tightness of two different pairs of upper/lower bounds for a general class of MIMO channels based on the unordered eigenvalue of the instantaneous correlation matrix and for arbitrary numbers of antennas. The tightness of Jensen’s bounds in different channel scenarios is investigated including multiuser MIMO with maximal ratio combining. Our analysis is facilitated by deriving some new results for finite-dimensional Wishart matrices, i.e., for the arbitrary moments of the unordered eigenvalue of central and non-central Wishart matrices. Our results provide very interesting insights into the implications of the system parameters, such as the number of antennas, and signal-to-noise ratio, on the tightness of Jensen’s bounds, and showcase the suitability and limitations of Jensen’s bounds.

Index Terms—Ergodic capacity, Jensen’s inequality, MIMO, random matrix theory, upper and lower bounds.

I. INTRODUCTION

The use of multiple antennas at both ends of a wireless link, known as multiple-input-multiple-output (MIMO) technology, can deliver substantial improvement in spectral efficiency and robustness, with no penalty in either power or bandwidth. For this reason, a vast amount of literature has been devoted to the analysis of the ergodic capacity of both single user (SU)-MIMO systems [1,2] and multiuser (MU)-MIMO systems [3,4].

Assuming a MIMO system with $N_t$ transmitters and $N_r$ receivers, the propagation channel is expressed by an $N_r \times N_t$ channel matrix $\mathbf{H}$, whose $(i,j)$th entry describes the path fading between the $j$th transmit and the $i$th receive antenna. Due to the intimate relationship between the performance of a MIMO system and the eigenvalue profile of the instantaneous correlation matrix $\mathbf{HH}^H$ (or $\mathbf{HH}^H$), the probabilistic characterization of these eigenvalues is necessary in order to derive analytical expressions for the ergodic capacity. For the prevalent Rayleigh and Ricean fading models, the instantaneous correlation matrix follows a complex Wishart (CW) distribution [5]. Fortunately, the eigenvalue statistics of CW matrices have been widely studied for many decades, for both the central [6,7] and non-central [7,8] cases. These results have been effectively applied to analyze the performance of SU-MIMO systems in different particular cases. More specifically, the joint probability density function (pdf) of ordered eigenvalues [9,10] was used to evaluate the information theoretical limits of SU-MIMO systems, while the pdfs of the smallest eigenvalue [11] and largest eigenvalue [12,13] have been utilized for designing antenna selection and multichannel beamforming schemes.

Although the performance of MIMO systems has been thoroughly investigated using random matrix theory tools, starting from independent and identically distributed (i.i.d.) Rayleigh fading [1,2,14,15] to the case where correlation is present at one of the two sides [8,11,16,17] or at both sides [18] or even Ricean fading channels [6,19]–[21], exact, closed-form capacity expressions are inherently hard to derive due to the complex distribution of the eigenvalues of $\mathbf{HH}$. In order to obtain more informative analytical insights into MIMO systems, a large number of works has focused on formulating upper and lower capacity bounds for both SU-MIMO (see [2,3,13,22]–[26]) and MU-MIMO systems [4,27] via Jensen’s inequality. Note that the Jensen’s inequality is a fundamental tool with a plethora of applications in areas such as probability, statistics [28,29] and communications.

However, the tightness of Jensen’s bounds (e.g. [23,25,30,31]) varies according to the eigenvalue profile of $\mathbf{HH}^H$, signal-to-noise ratio (SNR) and number of antennas. To the best of our knowledge, there appears to be a dearth of references investigating the tightness of Jensen’s bounds. A recent work [32], which leveraged the concentration of spectral measure of random matrices [33,34], applied a non-asymptotic analysis to characterize the power offset of Jensen’s bounds. The authors of [35] proposed an asymptotic analysis by expanding capacity

\footnote{The $(\cdot)^\dagger$ superscript denotes the Hermitian transpose.}

\footnote{Hereafter, we name the bounds on ergodic capacity, which are derived by applying Jensen’s inequality, as “Jensen’s bounds” for convenience.}
as an affine function adding a zero-order term, which is well suited for the high-SNR regime, while [36] gave a similar analysis in the low-SNR regime.

In this paper, we propose an alternative method to measure the tightness of Jensen’s bounds based on a basic mathematical idea—sandwich theorem [37]. By capitalizing on the sandwich theorem, we evaluate the offset between two pairs of upper and lower bounds by working exclusively with the eigenvalues of CW matrices. The offsets are analytically determined for several fading scenarios, i.e., i.i.d. Rayleigh, semi-correlated Rayleigh, Ricean channels for SU-MIMO systems and MU-MIMO with maximum ratio combining (MRC) receivers. Our results are based on the theory of finite-dimensional Wishart random matrices and, in particular, on the first positive and negative moment of the unordered eigenvalue of the central and non-negative and negative moment of the unordered eigenvalue of the Wishart random matrices and, in particular, on the first moments of the unordered eigenvalue of the Wishart random matrices. We evaluate the offset between two pairs of upper and lower bounds by working exclusively with the eigenvalues of the matrices and, in particular, on the first positive and negative moment of the unordered eigenvalue of the Wishart random matrices. We evaluate the offset between two pairs of upper and lower bounds by working exclusively with the eigenvalues of the matrices.

The channel $H_{su}$ can be effectively characterized by the matrix

$$H_{su} = \sqrt{\frac{K}{K+1}} M + \sqrt{\frac{1}{K+1}} R_{r}^{1/2} H_{d}^{1/2},$$

where $M$ is the deterministic channel component satisfying $\text{tr}(MM^H) = N_r N_t$, and $\tilde{H}$ is the random channel component containing i.i.d. $CN(0,1)$ entries. The Ricean $K$-factor is defined as the ratio of the power in $M$ and the average power in $\tilde{H}$. Moreover, $R_r \in \mathbb{C}^{N_r \times N_r} > 0$ and $R_d \in \mathbb{C}^{N_t \times N_t} > 0$ are the transmit and receive correlation matrices. In what follows, we refer to

$$m = \min(N_r,N_t), \quad n = \max(N_r,N_t)$$

and define the random matrix $\Phi \in \mathbb{C}^{m \times n}$ as

$$\Phi = \begin{cases} H_{su} H_{su}^*, & N_t \geq N_r, \\ H_{su} H_{su}^*, & N_t < N_r. \end{cases}$$

Also, let us denote, for convenience,

$$(R_m, R_n, M) = \begin{cases} (R_r, R_t, M), & N_t \geq N_r, \\ (R_t, R_r, M^H), & N_t < N_r. \end{cases}$$

Denoting by $\Lambda_m$ and $\Lambda_n$ the respective diagonal eigenvalue matrices of $R_m$ and $R_n$, the ergodic capacity can be written in the following two forms $^3$ [41, Eq. (2)]

$$C^{su} = E_{\tilde{H}} \left[ \log_2 \det \left( I_m + \frac{\lambda}{N_t} \Lambda_m \tilde{H} \Lambda_n \tilde{H}^H \right) \right],$$

and

$$C^{su} = m E_{\lambda} \left[ \log_2 \left( 1 + \frac{\lambda}{N_t} \right) \right].$$

$^3$When no CSI is available at the transmitter, the equal power allocation strategy, that appears in the expressions (6) and (7), is a robust transmission scheme since it maximizes the worst-case capacity in many cases [1,40] (i.e., the maxmin property). In the cases of general correlation structures and non-zero mean channels, this expression yields an achievable rate (lower bound on capacity) yet, we denote this expression as “capacity” for the sake of consistency with the vast body of MIMO literature [3,18,19,22,23,41,42].
where $\tilde{H} \sim CN_{m,n} \left( \sqrt{\frac{K}{K+1}} R_m^{-\frac{1}{2}} \tilde{M} R_n^{-\frac{1}{2}} \tilde{I}_n \otimes I_n \right)$, $\rho = P/\sigma^2$ is the average SNR, and $\lambda$ is the single unordered eigenvalue of $\Phi$. Assuming $f_{\lambda}(\Phi)$ is the marginal probability distribution function (pdf) of the unordered eigenvalue $\lambda$, we reformulate (7) as

$$C^{su} = m \int_0^\infty \log_2 \left( 1 + \frac{\rho}{\rho^t} \lambda \right) f_{\lambda}(\Phi) d\lambda. \quad (8)$$

This is the general capacity formula of SU-MIMO systems which encompasses many MIMO channels. However, in most cases of interest, it is really complicated, if not impossible, to evaluate analytically the theoretic expression in (8). Thus, for convenience, many prior works (see [22]–[24,41]) choose to characterize the ergodic capacity via its upper/lower bounds.

### III. General Analysis Framework for the Tightness of Jensen’s Bounds

In this section, we consider two pairs of Jensen’s bounds and quantify rigorously their offsets against the SNR. Note that these two pairs of bounds have been extensively applied in characterizing the capacity of MIMO systems (see [23,35]). Noting the fact that the evaluation of ergodic capacity can be reduced from the matrix problem in (6) into a real-valued scalar problem in (7), we first define two pairs of Jensen’s bounds for our general framework as

$$J_1 \triangleq \begin{cases} C^{su} \leq m \log_2 \left( 1 + \frac{\rho}{\rho^t} E \left[ \frac{1}{\lambda} \right] \right), \\ C^{su} \geq m \log_2 \left( 1 + \frac{\rho}{\rho^t} E \left[ \lambda^{-1} \right] \right), \end{cases} \quad (13)$$

and

$$J_2 \triangleq \begin{cases} C^{su} \leq E_{\Phi} \left[ \log_2 \left( \frac{\rho}{\rho^t} \det (\Phi) \right) \right] + m \log_2 \left( 1 + \frac{\rho}{\rho^t} E \left[ \lambda^{-1} \right] \right), \\ C^{su} \geq E_{\Phi} \left[ \log_2 \frac{\rho}{\rho^t} \det (\Phi) \right] + m \log_2 \left( 1 + \frac{\rho}{\rho^t} E \left[ \lambda^{-1} \right] \right). \end{cases} \quad (14)$$

It is important to note that $J_1$ and $J_2$ can be easily derived via the convexity or concavity of the corresponding functions, respectively. To evaluate the tightness of the pairs of bounds, we measure the offsets between the corresponding upper and lower bounds in $J_1$ and $J_2$.

**Theorem 1:** Let $\Phi$ be a $m \times m$ matrix as in (4) with unordered eigenvalue $\lambda$. The offsets of the two pairs of bounds defined in (13) and (14) are given by

$$\Delta_1 = m \log_2 \left( 1 + \frac{\omega - 1}{\rho \rho^t} E \left[ \frac{1}{\lambda} \right] + 1 \right) \quad (15)$$

and

$$\Delta_2 = m \log_2 \left( 1 + \frac{\omega - 1}{\rho \rho^t} E \left[ \lambda \right] + 1 \right) \quad (16)$$

respectively, where $\omega = E \left[ \frac{1}{\lambda} \right] E [\lambda]$.

**Proof:** The result can be derived directly by subtracting the lower bound from the upper bound in (13) and (14), respectively.

It is important to note that this result holds for arbitrary number of antennas and SNR. Since the offsets vary according to the statistical characteristics of the unordered eigenvalue, the tightness of both pairs of bounds is different for each MIMO channel. We now investigate the tightness of the proposed bounds as a function of the SNR. We compare the two offsets and determine which admits the smallest value depending on a SNR threshold.

**Corollary 1:** Define $\rho_t$ as the threshold such that

$$\begin{cases} \Delta_1 \leq \Delta_2, & \rho \leq \rho_t, \\ \Delta_1 > \Delta_2, & \rho > \rho_t, \end{cases} \quad (17)$$

SNR. Conditioned on $g_j$, $\tilde{g}_k$ is a Gaussian RV with zero mean and variance $\beta_k$ which does not depend on $g_j$. Therefore, $\tilde{g}_k$ is Gaussian distributed and independent of $g_j$ with $\tilde{g}_k \sim CN(0, \beta_k)$.
As the SNR grows large, i.e., \( \rho \to \infty \), the two offsets converge to
\[
\Delta_1 = m \log_2 (\omega), \quad \Delta_2 = 0. \tag{20}
\]
We see that the parameter \( N_t \) disappears in both offsets as \( \rho \) grows large, with \( \Delta_2 \) converging to zero, while \( \Delta_1 \) grows linearly with the minimum dimension of the channel matrix. In this case, the result shows that the capacity can be precisely predicted by \( J_2 \) at high-SNRs via the sandwich theorem.

As the SNR decreases, i.e., \( \rho \to 0 \), the two offsets converge to
\[
\Delta_2 = m \log_2 (\omega), \quad \Delta_1 = 0. \tag{21}
\]
Interestingly, we see that the two offsets have an inverse behavior in the low-SNR regime, which means that the properties of \( \Delta_2 \) at low-SNR can be derived by analyzing \( \Delta_1 \) in the high-SNR regime, and vice versa. In the latter case, the upper and lower bound in \( J_1 \) describe accurately the ergodic capacity.

Two insights may be useful:
\[
\Delta_{\rho_t} = \Delta_1 / 2, \quad \rho \to \infty. \tag{22}
\]
\[
\Delta_{\rho_t} = \Delta_2 / 2, \quad \rho \to 0. \tag{23}
\]
Since the maximum offset happens to be half the offset in the extreme SNR conditions, this symmetry of \( \Delta_1 \) and \( \Delta_2 \) is again illuminated.

To recap, Theorem 1 shows the symmetry of the two offsets around the SNR, which is very convenient for our following analysis since we can investigate the behavior of \( J_1 \) and \( J_2 \) in the high-SNR regime while omitting the low-SNR analysis.

IV. TIGHTNESS OF JENSEN’S BOUNDS IN FADING CHANNELS

Apart from the SNR, the tightness of Jensen’s bounds depends on the random channel matrix via its unordered eigenvalue. In this section, we also present some new results on Wishart random matrix theory, which are used to showcase the tightness of Jensen’s bounds for various fading models.

A. SU i.i.d. Rayleigh fading channels

For i.i.d. Rayleigh fading channels, the channel is statistically described by \( \tilde{H} \sim CN_{m,n} (0, I_m \otimes I_n) \). As shown in Theorem 1, to assess the tightness of \( \Delta_1 \) and \( \Delta_2 \), it is necessary to derive the first positive and first negative moment of the unordered eigenvalue of the channel matrix. Noting that the matrix \( \Phi \sim W_{m}(n, I_m) \), the desired results are given as [36, Lemma 4, Lemma 6]
\[
E [\lambda] = n, \tag{24}
\]
and
\[
E \left[ \frac{1}{\lambda} \right] = \frac{1}{n - m}, \quad n \neq m. \tag{25}
\]

Proposition 1: For i.i.d. Rayleigh fading channels, the two offsets of \( J_1 \) and \( J_2 \) with \( n \neq m \), are given in closed-form by
\[
\Delta_1 = m \log_2 \left( 1 + \frac{m \rho}{N_t + \rho (n - m)} \right) \tag{26}
\]
and
\[
\Delta_2 = m \log_2 \left( 1 + \frac{m N_t}{(\rho n + N_t) (n - m)} \right) \tag{27}
\]
respectively. Moreover, \( \rho_t \) is expressed as
\[
\rho_t = \frac{N_t}{\sqrt{n (n - m)}}. \tag{28}
\]

Proof: Substituting (24) and (25) into Theorem 1 we can obtain the desired result.

Our results in Proposition 1 give mathematical conclusions for the tightness of the two pairs of Jensen’s bounds, which are valid for all SNRs and arbitrary antenna configurations as long as \( n \neq m \). The results are insightful since they show explicitly that the two offsets reduce with increasing \( n \). A set of numerical results is presented in Fig. 1, where the theoretical results are compared against the empirical offsets, obtained via Monte-Carlo simulations. We see that the two offsets have indeed a mirror behavior with the symmetric center \( \rho_t \) varying with the dimensions of the channel matrix. Note that the Jensen’s bounds become inherently loose when the minimum number of antennas grows large (e.g., \( \Delta_1 \) grows up to 12.84 bits/Hz with \( m = 5 \)), since the variances of the involved random quantities become much higher. However, it is important to note that the two pairs of bounds are rather tight in the cases with small \( m \) and large \( n \), which correspond to a massive MIMO setup. To further illuminate the influence of matrix dimensions on the tightness of the bounds, we present the following corollaries.

Corollary 2: For i.i.d. Rayleigh fading channels, as the number of antennas on one side of the MIMO link grows large, i.e., \( n \to \infty \), \( \Delta_1 \) and \( \Delta_2 \) converge to 0.
The result holds for arbitrary m in the arbitrary SNR regime, and implies that both $J_1$ and $J_2$ are intimately tight for MIMO systems with large array antennas. In order to get more insights, we further investigate the tightness of $J_1$ for finite dimensions in the high-SNR regime.

**Corollary 3:** As the SNR grows large, i.e., $\rho \to \infty$, the offset of $J_1$ converges to

$$\Delta_1 = m \log_2 \left( \frac{n}{n-m} \right). \quad (29)$$

We see that, for high-SNR, the tightness of $J_1$ is only dependent on the numbers of antennas. The offset grows with the minimum number of antennas, while varying inversely with the maximum number of antennas. The result is confirmed in Fig. 2 as a function of $n$, for a system with $\text{SNR} = 20\, \text{dB}$. The minimum dimension of the channel matrices, which typically corresponds to the numbers of user antennas in practice, varies from 2 to 5. We find that all curves converge to zero with an increasing number of $n$, indicating that the Jensen’s bounds are particularly tight for massive MIMO configurations (e.g., $\Delta_1$ less than 0.5 bit/s/Hz with an antenna configuration $(4, 120)$).

Here, we further study $J_1$ by examining the following cases.

- For i.i.d. Rayleigh fading channels, adding $k$ antennas at both ends of link, $\Delta_1$ looses according to

$$\Delta_1 = (m + k) \log_2 \left( \frac{n + k}{n - m} \right). \quad (30)$$

- For i.i.d. Rayleigh fading channels, as the number of transmit and receive antennas grows with ratio $\beta = N_T/N_R \geq 1$, $\Delta_1$ simplifies to

$$\Delta_1 = m \log_2 \left( \frac{\beta}{\beta - 1} \right), \quad (31)$$

showing that the tightness of the bounds is inversely proportional to $m$ with fixed $\beta$. We point out that (30) and (31) showcase the impact of the minimum number of antennas, $m$, on the tightness of $J_1$ in the high-SNR regime and for arbitrary number of antennas.

### B. SU semi-correlated Rayleigh fading channels

For semi-correlated Rayleigh channels, we assume that the correlation is at the either end of the link, which makes the expressions for the marginal distribution of the unordered eigenvalue of $\Phi$ different. For convenience, we give a separate treatment of two cases, namely

$$(s,t, \Lambda_m, \Lambda_n) = \begin{cases} (m,n, \Theta_m, I_n), & \Theta_m \neq I_m, \\ (n,m, I_m, \Theta_n), & \Theta_n \neq I_n, \end{cases}$$

for the min semi-correlated and max semi-correlated Rayleigh fading channels, respectively. We now establish some new results on the unordered eigenvalue $\lambda$.

**Lemma 1:** For semi-correlated Rayleigh fading channels, with the unordered eigenvalue of $\Phi$ given as $\lambda$, the expectation of the arbitrary $p$th moment of $\lambda$ is given as

$$E[\lambda^p] = \frac{1}{m} \prod_{i<j}(\theta_j - \theta_i) \times \frac{\Gamma(t - s + j + p)}{\Gamma(t - s + j)} \det \left( \tilde{D}_j^p \right), \quad (32)$$

where $\nu = s - p + \text{sgn} (p)$, $s$, $\text{sgn} (\cdot)$ is the sign function, and $\tilde{D}_j^p$ is an $s \times s$ matrix whose $(l, k)$th entry is

$$\tilde{D}_j^p = \begin{cases} \theta_l^{k-1}, & l = 1, \ldots, s, \quad k \neq j, \\ \theta_l^{k+p-1}, & l = 1, \ldots, s, \quad k = j, \end{cases} \quad (33)$$

where $\theta_l$ is the $l$th eigenvalue of the correlation matrix $\Theta_m$ (or $\Theta_n$).

**Proof:** See Appendix A.

This lemma presents a new result for the statistical density of an unordered eigenvalue of a complex semi-correlated central Wishart matrix. It is important to mention that this
result holds for any \( p \), even for \( p < 0 \). From (32), we notice that when \( p < 0 \) with \( t < s \), some items in the expression become ill-defined. This is due to the fact that for max semi-correlated Rayleigh channels, the negative moments do not exist.

**Lemma 2:** For semi-correlated Rayleigh fading channels, the expectation of the unordered eigenvalue \( \lambda \) is given as

\[
E[\lambda] = \frac{t}{m} \sum_{i=1}^{s} \theta_i
\]

and its first negative moment for the min semi-correlated case is given by

\[
E\left[\frac{1}{\lambda}\right] = \frac{1}{m} \sum_{i=1}^{m} \frac{1}{\theta_i},
\]

while for the max semi-correlated case, \( E\left[\frac{1}{\lambda}\right] \) does not exist.

**Proof:** See Appendix B.

We note that (34) and (35) can be reduced to (24) and (25) when the correlation matrix is an identity matrix. Now, we are ready to derive the desired offsets for min semi-correlated Rayleigh fading channels.

**Proposition 2:** For min semi-correlated Rayleigh fading channels, the two offsets of \( J_1 \) and \( J_2 \) are given in closed-form as follows

\[
\Delta_1 = m \log_2 \left(1 + \frac{n \rho A - H \rho (n-m)}{N_1 + H \rho (n-m)}\right)
\]

and

\[
\Delta_2 = m \log_2 \left(1 + \frac{n N_1 A - H N_1 (n-m)}{(n \rho A H + H N_1) (n-m)}\right),
\]

where \( A = \frac{1}{m} \sum_{i=1}^{m} \theta_i \) represents the arithmetic mean of \( \theta_{i=1,...,m} \) and \( H = \frac{1}{m} \sum_{i=1}^{m} \frac{1}{\theta_i} \) is the harmonic mean of \( \theta_{i=1,...,m} \). Moreover, \( \rho_t \) is given by

\[
\rho_t = N_1 \sqrt{\frac{A}{(n-m)H}}.
\]

**Proof:** Substituting (34) and (35) into Theorem 1 and (18) yields the desired result.

It is important to note that the results give exact closed-form expressions for the offsets, which apply for all SNRs and arbitrary antenna configurations. Fig. 3 shows the offsets of two Jensen’s bounds for min semi-correlated Rayleigh channels. The correlation matrix \( \mathbf{R}_m \), whose \((i,j)\)th entry is \( e^{-0.05d^2(i-j)^2} \), corresponds to a d-wavelength antenna separation and a broadside Gaussian power azimuth spectrum with \( 2\pi \) root-mean-square spread [41]. The analytical expressions based on (36) and (37) match precisely with the simulation results, and perform quite similar as \( \Delta_1 \) and \( \Delta_2 \) in Fig. 1. By noting that the antenna configurations are identical to those considered in Fig. 1, it is straightforward to see that \( J_1 \) and \( J_2 \) are much looser in the semi-correlated channel scenario, since, in general, correlation increases the spread of the unordered eigenvalue.

**Corollary 4:** For min semi-correlated Rayleigh fading channels, as the SNR grows large, i.e., \( \rho \rightarrow \infty \), the offsets of \( J_1 \) and \( J_2 \) converge respectively to

\[
\Delta_1 = m \log_2 \left(\frac{n A}{(n-m) H}\right) \quad \text{and} \quad \Delta_2 = 0.
\]

**Proof:** This result can be directly derived by substituting (34) and (35) into (20).

As expected, the results in (39) reduce to (29) when the correlation matrix is an identity matrix. Therefore, the offsets in (39) show similar trend against the maximum number of antennas, as in the case of i.i.d. Rayleigh fading channels. It is important to note that both \( A \) and \( H \) depend on the \( m \) eigenvalues of the correlation matrix so that the impact of \( m \) cannot be separately analyzed.

**Corollary 5:** For min semi-correlated Rayleigh fading channels, as \( n \rightarrow \infty \), \( \Delta_1 \) reduces to

\[
\Delta_1 = m \log_2 \left(\frac{A}{H}\right) \geq 0 \quad \text{and} \quad \Delta_2 = 0.
\]

Note that the equality of \( \Delta_2 \) holds if and only if all elements in \( \Theta_m \) are the same.

From (40), we see that in a massive MIMO scenario, the tightness of \( J_1 \) is only related with the eigenvalues of the correlation matrix. We also note that the tightness of the bounds is influenced by the ratio of the arithmetic mean and harmonic mean of the eigenvalues of \( \Theta_m \). This ratio, as we known, is equal or greater than 1, and decreases with decreasing variance of \( \theta_i \). This implies that the Jensen’s bounds are less tight for semi-correlated Rayleigh channels compared with i.i.d. Rayleigh channels.

The results of Corollary 4 and Corollary 5 are confirmed in Fig. 4, where we compare the exact closed-form of offset \( \Delta_1 \) based on (36), and its asymptotic limit based on (40) for an SNR = 30dB. The correlation matrix is again modeled as
\[ \{ R_{m} \}_{i,j} = e^{-0.2(i-j)^2}. \]

In all cases, given as \( m = 2, 3, 4 \), there is an exact agreement between the analytical results and the Monte-Carlo simulations. However, we see that the asymptotic limit of \( \Delta_1 \) is larger than 6.5 bits/s/Hz when \( m = 4 \); this means that \( J_1 \) is not suitable to measure the ergodic capacity in the high SNR regime when the number of user antennas grows large, even for a very large \( n \). As shown in Fig. 3, \( J_2 \) would be more appropriate for semi-correlated Rayleigh channels when the SNR grows large.

\[ \begin{array}{ccc}
1 & 2 & 3 \\
4 & 5 & 6 \\
7 & 8 & 9 \\
10 & Analytical in (36) \\
Monte-Carlo \\
Asymptotic: \Delta_1 \text{ in (40)} \\
\end{array} \]

\( \Delta_1 \) for min semi-correlated Rayleigh channels against the maximum number of antennas with \( m = 2, 3, 4 \), and \( \rho = 30 \text{dB} \).

C. **SU i.i.d. Ricean fading channels**

For i.i.d. Ricean fading channels, we model \( \hat{H} \sim CN_{m,n} \left( \sqrt{\frac{k}{K+1}} M, \frac{1}{K+1} I_m \otimes I_n \right) \) by simplifying \( \Lambda_m = I_m \) and \( \Lambda_n = I_n \). With this model, \( \Phi \) follows a complex noncentral Wishart distribution, denoted as \( \Phi \sim W_m \left( n, \frac{1}{K+1} I_m, \Xi \right) \), where \( \Xi \equiv K \Xi M M^\dagger \).

**Lemma 3:** Let \( \epsilon_1, \ldots, \epsilon_m \) be the \( m \) eigenvalues of \( \Xi \). Then, the first positive and first negative moment of the singular unordered eigenvalue \( \lambda \) of \( \Phi \) are

\[ E[\lambda] = \frac{\kappa}{2(K+1)} \det \{ U \} \]  
(41)

and

\[ E \left[ \frac{1}{\lambda} \right] = 2\kappa (K+1) \det \{ \hat{U} \}, \]  
(42)

respectively, where \( \kappa \equiv \sqrt{2^{m-k}m^{(m+1)/2}} \). Note that \( U \) and \( \hat{U} \) are \( m \times m \) matrices whose \((i, k)\)th entry is given by

\[ \{ U \}_{i,k} = \begin{cases} Q_{n-m+2k-1,n-m} \left( \sqrt{2\epsilon_i}, 0 \right), & k = 1, \ldots, m-1, \\ Q_{n+m+1,n-m} \left( \sqrt{2\epsilon_i}, 0 \right), & k = m. \end{cases} \]  
(43)

and

\[ \{ \hat{U} \}_{i,k} = \begin{cases} Q_{n-m-1,n-m} \left( \sqrt{2\epsilon_i}, 0 \right), & k = 1, \\ Q_{n-m+2k-1,n-m} \left( \sqrt{2\epsilon_i}, 0 \right), & k = 2, \ldots, m. \end{cases} \]  
(44)

respectively. In the above, \( Q_{p,q}(a,b) \) is the Nutall \( Q \)-function, defined in [21] as

\[ Q_{p,q}(a,b) = \int_b^\infty x^p e^{-\frac{x^2+a^2}{2}} I_q(a x) \, dx, \]  
(45)

and \( I_q(\cdot) \) is the \( q \)th order modified Bessel function of the first kind.

**Proof:** See Appendix C.

We see that the results in (41) and (42) can be easily programmed and efficiently calculated. Additionally, the Nutall \( Q \)-functions in (43) and (44) have closed-form representations shown in (93) since the lower limit of the integral in (45) is equal to 0.

**Proposition 3:** For i.i.d. Ricean fading channels, the two offsets of \( J_1 \) and \( J_2 \) are given in closed-form by

\[ \Delta_1 = m \log_2 \left( 1 + \frac{\rho \kappa^2 \det \{ \hat{U} \} \det \{ U \} - \rho}{\kappa \det \{ U \} + \rho} \right), \]  
(46)

and

\[ \Delta_2 = m \log_2 \left( 1 + \frac{\gamma \kappa^2 \det \{ \hat{U} \} \det \{ U \} - \gamma}{\rho \kappa \det \{ U \} + \gamma} \right), \]  
(47)

respectively, where \( \gamma \equiv 2N_l (K+1) \). Moreover, \( \rho_i \) is given by

\[ \rho_i = 2 (K+1) N_l \sqrt{\frac{\det \{ \hat{U} \}}{\det \{ U \}}}. \]  
(48)

**Proof:** Substituting (41) and (42) into Theorem 1 and (18) yields the desired result.

It is important to note that the results in Proposition 3 apply for arbitrary dimensions of the channel matrix and all SNRs. However, it is hard to unveil the impact of each parameter on the offsets due to the presence of the determinant inside the logarithmic expression. In order to derive more informative insights, we analyze the single-antenna user case, and obtain a simpler result given in the next corollary.

**Corollary 6:** For i.i.d. Ricean fading systems with a single-antenna user, i.e., \( m = 1 \), the closed-form expressions of the two offsets are given as

\[ \Delta_1 = \log_2 \left( 1 + \frac{2\rho \epsilon (n+\epsilon) - 2\rho}{\gamma \epsilon + 2\rho} \right), \]  
(49)

and

\[ \Delta_2 = \log_2 \left( 1 + \frac{\gamma \epsilon (n+\epsilon) - \gamma}{2\rho(n+\epsilon) + \gamma} \right), \]  
(50)

respectively, where \( \epsilon \equiv 2N_l (K+1) \) and \( \epsilon \) is the eigenvalue of \( \Xi \in \mathbb{C}^{1 \times 1} \).

**Proof:** See Appendix D.

The above results place no restrictions on the maximum dimension of the matrix. Most importantly, the two offsets for single input multiple output systems can be easily evaluated since they only involve standard functions.

**Corollary 7:** For i.i.d. Ricean fading systems with single-
antenna users, as $n \to \infty$,
\[
\lim_{n \to \infty} \Delta_1 = 0. \tag{51}
\]

**Proof:** By applying a Taylor series expansion, $e^x$ can be divided into two parts, given as
\[
e^{-x} = \sum_{k=0}^{n-2} \frac{(-e)^k}{k!} + \sum_{k=n-1}^{\infty} \frac{(-e)^k}{k!}. \tag{52}
\]
By noting that $\xi$ involves the first part on the right side of (52), we substitute (52) in $\xi$ and after some basic manipulations, we obtain
\[
\xi = O \left( \frac{1}{n-1} \right), \tag{53}
\]
where $O(\cdot)$ is expressed as $f(n) = O(g(n))$ as $n \to \infty$, if $\lim_{n \to \infty} \frac{f(n)}{g(n)} = 1$. Then, substituting (53) into (49) and (50) yields the result. \hfill \Box

**D. MU-MIMO MRC Systems**

For a MU-MIMO MRC system modeled as in (11), whose spectral efficiency is given in (12), we note that the desired signal power and interference power are independent of each other, which makes it possible to derive them separately. We first notice that $\{g_jH\} \sim W_1(M, I_M)$ in (12), thus, we simply get [36, Lemma 4]
\[
E \left[ \rho_{mu} g_jH \right] = \rho_{mu} M \beta_j \tag{54}
\]
and
\[
E \left[ \left( \rho_{mu} g_jH \right)^{-1} \right] = \frac{1}{\rho_{mu} (M-1) \beta_j}. \tag{55}
\]
Then, an exact expression of the expectation of the noise-plus-interference (NPI) power is derived, as well as an approximate formula of the first negative moment of the NPI power.

**Lemma 4:** The first positive and first negative moment of the noise-plus-interference power $I = 1 + \rho_{mu} \sum_{k=1, k \neq j}^{N} |\tilde{g}_k|^2$ are given as
\[
E[I] = 1 + \rho_{mu} \varpi \tag{56}
\]
and
\[
E[I^{-1}] = \sum_{k=1, k \neq j}^{N} \frac{-L(\rho_{mu} \beta_k)}{\rho_{mu} T(j, k, B)}, \tag{57}
\]
where $L(x) = e^{\frac{x}{2}} \text{Ei} \left( -\frac{1}{x} \right)$, $T(j, k, B) = \beta_k \prod_{l=1, l \neq k, l \neq j}^{N} (1 - \beta_l / \beta_k)$, $\varpi = \sum_{k=1, k \neq j}^{N} \beta_k$ and $\text{Ei} (\cdot)$ is the exponential integral function [38, Eq. (5.1.2)].

**Proof:** See Appendix E. \hfill \Box

Note that the results in (56) can be very easily evaluated. Both $E[I]$ and $E[I^{-1}]$ are only related to the users’ positions while independent of the number of antennas at the BS. This is confirmed in Fig. 6, in which we compare the analytical expression of $E[I^{-1}]$ based on (57) and Monte-Carlo results for a MU-MIMO MRC system. We set the number of antennas at BS as $M = 128$, and assume that the number of users $N = 2, 5, 15$. All users are dropped randomly in a cell with radius of 1000 meters, and we assume that no user is closer to the BS than $r_h = 100$ meters. Without loss of generality, the large-scale fading is modeled via $\beta_j = z_j / (r_j / r_h)^\nu$, where $z_j$ obeys a log-normal distribution with standard deviation $\sigma_{\text{shadow}}$ and $r_j$ is the distance from the $j$th user to BS. In all cases, the curves remain constant at 1 at low SNRs, and undergo a sharp decrease until decreasing to 0 at high SNR. Moreover, there appears to be a significant property that the number of users has an profound impact on $E[I^{-1}]$; in particular, $E[I^{-1}]$ decreases with increasing $N$ since interference $I$ increases linearly with the number of users.

**Lemma 5:** For the function as $L(x)$, the following properties are satisfied:
\[
\lim_{x \to \infty} \frac{L(x)}{\ln(x)} = -1, \quad \text{and} \quad \lim_{x \to 0} \frac{L(x)}{x} = -1. \tag{58}
\]

**Proof:** The results can be obtained by using L’Hôpital’s Rule directly. \hfill \Box

**Proposition 4:** For a MU-MIMO system with an MRC receiver, the expressions of the two offsets of $J_1$ and $J_2$ are
given by
\[
\Delta_1 = \log_2 \left( 1 + M \beta_j \sum_{k=1,k \neq j}^N \frac{-L(\rho_{mu} \beta_k)}{T(j,k,B)} \right) - \log_2 \left( 1 + \frac{\rho_{mu} (M-1) \beta_j}{1 + \rho_{mu} \omega} \right)
\]
and
\[
\Delta_2 = \log_2 \left( 1 + \frac{\rho_{mu} \omega + 1}{\rho_{mu} (M-1) \beta_j} \right) - \log_2 \left( 1 + \frac{1}{M \beta_j \sum_{k=1,k \neq j}^N \frac{-L(\rho_{mu} \beta_k)}{T(j,k,B)}} \right).
\]
(59)

**Proof**: Substituting (56) and (57) into Theorem 1 yields the desired result.

This result shows the relationship between the large-scale fading coefficient and tightness of bounds, which varies with the desired result.

\[
\log_2 \left( 1 + \frac{M \beta_j}{1 + \rho_{mu} \omega} \right)
\]
and
\[
\log_2 \left( 1 + \frac{1}{M \beta_j \sum_{k=1,k \neq j}^N \frac{-L(\rho_{mu} \beta_k)}{T(j,k,B)}} \right).
\]
(60)

- As \(\rho_{mu} \to 0\), the two offsets reduce to

\[
\Delta_2 = \log_2 \left( \frac{M}{M-1} \sum_{k=1,k \neq j}^N \frac{\beta_k}{T(j,k,B)} \right), \text{ and } \Delta_1 = 0
\]

The results are obtained by utilizing Lemma 5 in (59) and (60). We see that \(J_1\) can perfectly characterize the capacity in the low-SNR regime. Moreover, the tightness of \(J_2\) depends on the LSF from each user, and when \(M\) grows large, \(\Delta_2\) converges to

\[
\Delta_2 = \log_2 \left( \frac{N}{\sum_{k=1,k \neq j}^N \frac{\beta_k}{T(j,k,B)}} \right).
\]
(62)

- As \(\rho_{mu} \to \infty\), the two offsets approximately reduce to

\[
\Delta_1 = \log_2 \left( \frac{\omega + \omega M \beta_j \sum_{k=1,k \neq j}^N \ln \beta_k}{(M-1) \beta_j} \right)
\]

and

\[
\Delta_2 = \log_2 \left( 1 + \frac{\omega}{(M-1) \beta_j} \right) - \log_2 \left( 1 + \frac{1}{M \beta_j \sum_{k=1,k \neq j}^N \frac{\ln \beta_k}{T(j,k,B)}} \right).
\]
(63)

This result is obtained by using Lemma 5 and [43, Lemma 2] in Proposition 4. The result shows the significant impact of LSF on the tightness of bounds in the high-SNR regime with a finite number of antennas.

- In the massive MIMO regime, i.e., as \(M \to \infty\), the two offsets approximately converge to

\[
\Delta_1 = \log_2 \left( 1 + \frac{\rho_{mu} \omega}{\rho_{mu}} \sum_{k=1,k \neq j}^N \frac{-L(\rho_{mu} \beta_k)}{T(j,k,B)} \right)
\]

and

\[
\Delta_2 = 0.
\]
(65)

Note that \(J_2\) is considerably tight for massive MIMO for any the arbitrary SNRs. Moreover, we see that as the SNR grows large, \(\Delta_1\) further converges to

\[
\Delta_1 = \log_2 \left( \frac{\omega}{\sum_{k=1,k \neq j}^N \frac{\ln \beta_k}{T(j,k,B)}} \right).
\]
(66)

We note that the LSF of each user has significant influence on the tightness of Jensen’s bounds. Therefore, in contrast of SU-MIMO scenarios, the two offsets do not converge to the same value in the high and low SNR regime for MU-MIMO MRC systems. However, from (65), we see that \(J_2\) is rather tight for arbitrary SNRs in massive MIMO systems.
In Fig. 7, the offsets of $J_1$ based on (59) and $J_2$ based on (60) are depicted with 10 users in the cell. Here, we consider the same parameters as before. We first note that the absolute values of offsets of both bounds are rather small, which implies that the tightness of $J_1$ and $J_2$ is confirmed to be very satisfactory with single-antenna users. Moreover, we see that $\Delta_1$ does not have a mirror behavior against $\Delta_2$ in contrast with SU-MIMO systems, i.e., both bounds are tighter in the low-SNR regime. However, it still holds that $J_1$ performs better at low-SNRs while $J_2$ is tighter at high-SNRs. Similar to the SU cases, the tightness of bounds also improves with increasing the maximum number of antennas.

V. DISCUSSION

In this section, we further investigate the analytical offsets of Jensen’s bounds for the SU-MIMO and MU-MIMO systems, and present a method of minimizing the offsets by selecting proper bounds in different SNR regimes. Note that all the results in this section are derived by prior expressions in Section IV.

In order to show the implications of our results, we analyze the most common scenarios of interest shown in Table I. We assume that the system works at $-10$dB and $15$dB corresponding to low-SNR and high-SNR regimes, respectively. For single user cases, we consider that the user is equipped with 1, 2 or 4 antennas while the base station (BS) has 2, 4, 8 or 128 antennas (i.e., $m = 1, 2, 4$ and $n = 2, 4, 8, 128$) to represent the finite-dimension MIMO systems and massive MIMO systems. Note that in i.i.d Ricean fading channels, the deterministic component $M$ with singular values $\sigma_k = 1, ..., m$ is randomly generated shown in Table II. For multiuser cases, we only consider there are 2 or 4 single-antenna users (i.e., $m = 2, 4$) in the system. The deterministic LSF $\beta_k$ is derived based on the location of $k$th user who is randomly distributed in the area, as shown in Table II.

For SU-MIMO, from Table I, it can be confirmed that $J_1$ is more suitable for low-SNR scenarios while $J_2$ performs better at high-SNRs. The offset $\Delta_2$ is rather small, numerically, less than 0.5bit/s/Hz in the most of cases when the system is working at $15$dB regardless of the fading channels, especially for the cases of small number of terminal antennas. Note that the offset can be further reduced as $\rho$ grows based on our prior analysis in Section III. In other words, the numerical results indicate that the capacity can be approximately expressed by both the upper and lower bound in $J_2$ with the offset less than $\Delta_2$ at $15$dB (i.e., less than $0.087$ bits/Hz for $m \times n = 2 \times 4$). On the contrary, $J_1$ behaves poorly in this regime while it becomes rather tight at $-10$dB. Therefore, we can infer that the capacity can be simply measured by selecting the right pairs of Jensen’s bounds in different SNR regimes without calculating it, and show how tight the bounds are. We also note that Jensen’s bounds perform even tighter with 128 antennas at the BS as expected, which validates that the Jensen’s bounds become tighter for massive MIMO.

For MU-MIMO, recall that the tightness of Jensen's bounds is effected by the location of users. Without loss of generality, we randomly generate the users and model LSF via $\beta_k = z_k/(r_k/r_h)$ shown in Table II. We first note that the Jensen’s bounds can be used to measure the capacity in the low-SNR regimes since both pairs of them behave really tight at $-10$dB. However, we are more interested in the behavior of Jensen’s bounds in the high-SNR regime. The numerical results show that the Jensen’s bound is not as tight as in the cases of low-SNRs, but still acceptable in representing capacity especially in the massive MIMO scenarios.

So far, we have confirmed that the capacity can be precisely predicted by proper Jensen’s bound in both low-SNR and high-SNR regimes. We now investigate the worst performance of Jensen’s bounds for SU-MIMO with respect to SNR. As shown in Corollary 1, the offset $\Delta_{\rho_1}$ at $\rho_1$ can be regarded as the maximum offset since we can select different Jensen’s bound.

### Table I

<table>
<thead>
<tr>
<th>$\rho$ [dB]</th>
<th>$m \times n$</th>
<th>Single user</th>
<th>Multiuser</th>
</tr>
</thead>
<tbody>
<tr>
<td></td>
<td>i.i.d. Rayleigh</td>
<td>semi-correlated Rayleigh</td>
<td>i.i.d. Ricean</td>
</tr>
<tr>
<td></td>
<td>$\Delta_1$</td>
<td>$\Delta_2$</td>
<td>$\Delta_1$</td>
</tr>
<tr>
<td></td>
<td>$\Delta_1$</td>
<td>$\Delta_2$</td>
<td>$\Delta_1$</td>
</tr>
<tr>
<td>-10</td>
<td>1 x 2</td>
<td>0.067</td>
<td>0.933</td>
</tr>
<tr>
<td></td>
<td>2 x 4</td>
<td>0.134</td>
<td>1.866</td>
</tr>
<tr>
<td></td>
<td>4 x 8</td>
<td>0.269</td>
<td>1.372</td>
</tr>
<tr>
<td></td>
<td>2 x 128</td>
<td>0.005</td>
<td>0.053</td>
</tr>
<tr>
<td></td>
<td>4 x 128</td>
<td>0.021</td>
<td>0.214</td>
</tr>
</tbody>
</table>

### Table II

<table>
<thead>
<tr>
<th>$m \times n$</th>
<th>Single user i.i.d. Ricean</th>
<th>Multiuser i.i.d. Rayleigh</th>
</tr>
</thead>
<tbody>
<tr>
<td></td>
<td>${k_{1}=1,...,m}$</td>
<td>${\beta_{k=1,...,m}}$</td>
</tr>
<tr>
<td>1 x 2</td>
<td>{2}</td>
<td>{2}</td>
</tr>
<tr>
<td>2 x 4</td>
<td>{0.343, 7.617}</td>
<td>{0.343, 7.617}</td>
</tr>
<tr>
<td>4 x 8</td>
<td>{1.463, 2.358, 3.519, 24.66}</td>
<td>{0.051, 0.150, 0.056, 0.139}</td>
</tr>
<tr>
<td>2 x 128</td>
<td>{30.44, 225.6}</td>
<td>{0.065, 0.218}</td>
</tr>
<tr>
<td>4 x 128</td>
<td>{24.12, 31.09, 40.99, 415.8}</td>
<td>{0.051, 0.150, 0.056, 0.139}</td>
</tr>
</tbody>
</table>
to reduce the offset at other SNR values. Thus, we evaluate both Δρi and ρt shown in Table III with the same simulation parameters shown in Table II.

For the cases of m = 1, we note two features: (a) ρi admits values which are quite smaller than practical values of interest; (b) the maximum offset Δρi is rather small. The results indicate that Jensen’s bounds are reliable in characterizing the capacity of SIMO with arbitrary fading channels. For the cases of m = 2, 4, we see that the offsets become larger, especially when the user is equipped with 4 antennas. Therefore, Jensen’s bounds can hardly predict the capacity around ρi when multi-antennas are deployed at terminal side. However, it is important to note that ρi is still much less than the usual working SNR of the concurrent practical systems.

VI. CONCLUSION

This paper has investigated a problem in the context of MIMO systems that has been surprisingly overlooked in the relevant literature, namely the tightness of Jensen’s capacity bounds. In order to perform this task, we have first derived some new results for finite-dimensional Wishart random matrix theory, by elaborating on central and non-central Wishart matrices with arbitrary dimensionality. We point out that our results for the unordered eigenvalue of different classes of Wishart matrices are closed-form expressions that follow:

\[ E[λ^p] = \frac{1}{m!} \prod_{i<j} (θ_j - θ_i) \times \sum_{i=1}^{m} \sum_{j=1}^{m} \frac{Γ(n-m+j+p)}{Γ(n-m+j)} \frac{θ_j^{-i}θ_i^{m-n-1}}{D_{i,j}}. \]  

(69)

The \( (s, A_s, A_n) = (I_s, l_s, Θ_n) \) Case: For this case, we start by employing a result from [23, Eq. (14)] to express the pdf of the unordered eigenvalue θi, conditioned on Θn, as follows:

\[ f_λ(λ) = \frac{1}{m!} \prod_{i<j} (θ_j - θ_i) \times \sum_{i=1}^{n} \sum_{j=1}^{n} \frac{Γ(n-m+j+p)}{Γ(n-m+j)} \frac{θ_j^{-i}θ_i^{m-n-1}}{D_{i,j}}. \]  

(70)

The arbitrary order moment of λ is calculated in the same way as in (69).

By setting s equal to the number of antennas on the side with correlation and t equal to the number of antennas on the
other side, we obtain
\[
E [\lambda^p] = \frac{1}{m \prod_{i<j}^s (\theta_j - \theta_i)} \sum_{j=s-m+1}^{s} \frac{\Gamma (t-s+j+p)}{\Gamma (t-s+j)} \det \left( \tilde{D}_j^p \right). \tag{71}
\]
For the \( p > 0 \) case, \( \det (\tilde{D}_j^p) = 0 \) when \( j \leq s-p \), while for the \( p < 0 \) case, \( \det (\tilde{D}_j^p) = 0 \) when \( j > -p \). By noting this property, we conclude the proof.

**APPENDIX B**

**PROOF OF LEMMA 2**

In order to prove this corollary, we first analyze the properties of the generalized Vandemonde determinants.

**Lemma 6:** Let us define a generalized Vandemonde matrix according to

\[
V_{n-k} = \begin{pmatrix}
1 & 1 & 1 & \cdots & 1 \\
\alpha_1 & \alpha_2 & \alpha_3 & \cdots & \alpha_n \\
\vdots & \vdots & \vdots & \cdots & \vdots \\
\alpha_{n-k+1}^{n-k-1} & \alpha_{n-k+2}^{n-k-1} & \alpha_{n-k+3}^{n-k-1} & \cdots & \alpha_n^{n-k-1} \\
\vdots & \vdots & \vdots & \cdots & \vdots \\
\alpha_1^k & \alpha_2^k & \alpha_3^k & \cdots & \alpha_n^k 
\end{pmatrix}, \tag{72}
\]

where \( k = 0, 1, \ldots, n \). The determinant of \( V_{n-k} \) is equal to

\[
\det (V_{n-k}) = \sum_{q \geq \cdots \geq j > i} \alpha_i \cdots \alpha_q \prod_{i<j}^n (\alpha_j - \alpha_i). \tag{73}
\]

**Proof:** We first extend (72) by adding one column \( y^j \) and one row \( \alpha_i^{n-k} \) with \( i, j = 0, 1, \ldots, n \), as follows

\[
V = \begin{pmatrix}
1 & 1 & 1 & \cdots & 1 \\
\alpha_1 & \alpha_2 & \alpha_3 & \cdots & \alpha_n & y \\
\vdots & \vdots & \vdots & \cdots & \vdots & \vdots \\
\alpha_{n-k+1}^{n-k-2} & \alpha_{n-k+2}^{n-k-2} & \alpha_{n-k+3}^{n-k-2} & \cdots & \alpha_n^{n-k-2} & y^{n-k} \\
\alpha_1^k & \alpha_2^k & \alpha_3^k & \cdots & \alpha_n^k & y^k 
\end{pmatrix}. \tag{74}
\]

Recalling the property of Vandemonde determinants, we have

\[
\det (V) = \prod_{i=1}^n (y - \alpha_i) \prod_{i<j}^n (\alpha_j - \alpha_i) \\
= \left( y^n + (-1)^{n-1} \sum_{i=1}^n \alpha_i + \cdots \right) \\
+ (-1)^k y^{n-k} \sum_{q > \cdots > j > i} \alpha_i \alpha_j \cdots \alpha_q + \cdots \\
+ (-1)^n \prod_{i=1}^n \alpha_i \prod_{i<j} (\alpha_j - \alpha_i). \tag{75}
\]

We also notice that \( \det (V) \) can be calculated by expanding it along its last row as

\[
\det (V) = \sum_{i=0}^n (-1)^{n-i} y^{n-i} \det (V_{n-i}). \tag{76}
\]

Comparing (75) and (76) in terms of the factor \( y^k \), we obtain the result.

By establishing Lemma 6, the expectation of \( \lambda \) can be simplified as

\[
E [\lambda] = \frac{t}{m} \sum_{i=1}^s \theta_i. \tag{77}
\]

To derive \( E \left[ \frac{1}{\lambda} \right] \), we generate a generalized Vandemonde matrix as

\[
U = \begin{pmatrix}
\frac{1}{\pi_1} & \frac{1}{\pi_2} & \frac{1}{\pi_3} & \cdots & \frac{1}{\pi_n} \\
1 & 1 & 1 & \cdots & 1 \\
\vdots & \vdots & \vdots & \cdots & \vdots \\
\frac{1}{\pi_1^{n-2}} & \frac{1}{\pi_2^{n-2}} & \frac{1}{\pi_3^{n-2}} & \cdots & \frac{1}{\pi_n^{n-2}} \\
\frac{1}{\pi_1^{n-1}} & \frac{1}{\pi_2^{n-1}} & \frac{1}{\pi_3^{n-1}} & \cdots & \frac{1}{\pi_n^{n-1}} \\
\frac{1}{\pi_1^k} & \frac{1}{\pi_2^k} & \frac{1}{\pi_3^k} & \cdots & \frac{1}{\pi_n^k}
\end{pmatrix}, \tag{78}
\]

whose determinant can be easily obtained as

\[
\det (U) = \prod_{i=1}^n \frac{(\pi_j - \pi_i)}{\xi_i}. \tag{79}
\]

Then, we calculate \( \det (U_{n-k}) \) in the same way as in (73). Thus, we get

\[
\det (U_{n-k}) = \sum_{q > \cdots > j > i} \alpha_i \alpha_j \cdots \alpha_q \prod_{i<j}^{n-k} (\alpha_j - \alpha_i) / \prod_{i=1}^n \alpha_i. \tag{80}
\]

We notice the fact that the expectation of \( 1/\lambda \) can also be calculated as

\[
E \left[ \frac{1}{\lambda} \right] = \frac{\Gamma (t-s) \det \left( \tilde{D}_1^{-1} \right)}{m \Gamma (t-s + 1) \prod_{i<j} (\theta_j - \theta_i)}, \tag{81}
\]

which has a zero mass, only if \( t - s > 0 \). Thus, the first negative moment of \( \lambda \) does not exist when correlation is on the side with the maximum number of antennas, while

\[
E \left[ \frac{1}{\lambda} \right] = \frac{1}{m (n-m)} \sum_{i=1}^s \frac{1}{\theta_i}. \tag{82}
\]

when correlation is on the side with the minimum number of antennas.

**APPENDIX C**

**PROOF OF LEMMA 3**

We first give some preliminary results.

**Lemma 7:** Let \( \epsilon_1, \ldots, \epsilon_m \) be the \( m \) eigenvalues of \( \Xi \). The marginal density distribution of a single eigenvalue \( \lambda \) of \( \Phi \sim \mathcal{W}_m \left( \frac{1}{n+k+1}, I_m, \Xi \right) \) is given by (83) shown at the top of the next page, where \( U_{i,j} \) is the \( (i,j) \)-cofactor of an \( m \times m \) matrix whose \( (l,k) \)-th entry is

\[
\{ U \}_{l,k} = (n-m)_{l+1}^{n-m} \frac{m-n+2k+2}{2} e^{\epsilon_l} Q_{n-m+2k-1,n-m} (\sqrt{2 \epsilon_l}/0). \tag{84}
\]

**Proof:** The joint pdf of the unordered eigenvalues of \( \Phi \sim \mathcal{W}_m \left( \frac{1}{n}, I_m, \Xi \right) \) has been given in [19, Theorem 1]. After some basic transformations, we derive the joint pdf of the eigenvalues \( \lambda_1, \ldots, \lambda_m \) of \( \Phi \), given by (85). To evaluate the
\[ f_\lambda (\lambda) = \frac{\text{etr} (\Xi)}{m((n-m)!)^m \prod_{k<i} (\epsilon_i - \epsilon_k)} \times  \sum_{i=1}^{m} \sum_{j=1}^{m} e^{-(K+1)\lambda} \left((K+1)\lambda\right)^{n-m+j} {}_0F_1 (n-m+1, (K+1)\epsilon_i\lambda) U_{i,j} \]  

(83)

\[ f_\lambda (\lambda_1, \ldots, \lambda_m) = \frac{\text{etr} (\Xi)}{((n-m)!)^m \prod_{k<i} (\epsilon_i - \epsilon_k)} \times \det \{0 F_1 (n-m+1, (K+1)\epsilon_i\lambda)\} \det \left\{((K+1)\lambda_i)^{j-1}\right\} \prod_{i=1}^{m} e^{-(K+1)\lambda_i} \lambda_i^{n-m} \frac{1}{(K+1)^{m-n+1}} \]  

(85)

\[ \{U\}_{i,k} = (K+1)^{n-m+k} \int_0^\infty e^{-(K+1)\lambda\lambda^{n-m+k}} {}_0F_1 (n-m+1, (K+1)\epsilon_i\lambda) d\lambda \]  

(86)

pdf of the unordered single eigenvalue of \( \Psi \), we apply [42, Lemma 2] in (85). It is then straightforward to get (83). The \((l, k)\)-th entry of \( U \) can be calculated in the same way as in [21, Eq. (60)] and is given in (86).

With Lemma 7, using the linearity of the determinants and after some basic manipulations, the expectation of \( \lambda \) can be derived as (87) shown at the top of the next page, where \( U_j \) is \( m \times m \) matrix whose \((l, k)\)-th entry is

\[ \{U\}_{j,k} = \begin{cases} Q_{n-m+2k-1-n-m} (\sqrt{2\epsilon_i} 0), & k \neq j, \\ Q_{n-m+2k+1-n-m} (\sqrt{2\epsilon_i} 0), & k = j. \end{cases} \]  

(88)

Noting that \( \det (U_j) = 0 \) when \( j \leq m - 1 \), we obtain the result. Note that the expectation of \( 1/\lambda \) can be derived in the same way as in (87).

\section*{APPENDIX D}

\section*{PROOF OF COROLLARY 6}

Letting \( m = 1 \), it is straightforward to obtain

\[ \kappa = 2 \frac{\lambda}{\epsilon} e^{\frac{\lambda}{2\epsilon}}, \]  

(89)

\[ E[\lambda] = Q_{n-2-n-1} \left( \sqrt{2\epsilon} 0 \right), \]  

(90)

and

\[ E\left[ \frac{1}{\lambda} \right] = Q_{n+2-n-1} \left( \sqrt{2\epsilon} 0 \right). \]  

(91)

We notice that \( I_q (\cdot) \) in (45) can be transformed into a first kind Bessel function [44, Eq. (8.406.3)] as

\[ I_v (z) = i^{-v} J_v \left( iz \right), \]  

(92)

where \( i = \sqrt{-1} \). Substituting (92) in (45) and using [44, Eq. (6.631.1)], \( Q_{p,q} (a, 0) \) can be derived as

\[ Q_{p,q} (a, 0) = i^{-q} \int_0^\infty x^p e^{-\frac{x^2}{2}} J_q \left( iax \right) dx \]  

\[ = i^{-q} e^{-\frac{a^2}{2}} \int_0^\infty x^p e^{-\frac{x^2}{2}} J_q \left( iax \right) dx \]  

\[ = a^q \Gamma \left( \frac{p+q+1}{2} \right) F_1 \left( \frac{p+q+1}{2}; q+1; \frac{a^2}{2} \right). \]  

(93)

We can now simplify the expectation of \( \lambda \) as

\[ E[\lambda] = \frac{\left( \sqrt{2\epsilon} \right)^{n-1}}{2(n-1)} \Gamma (n-1) \left( 1 - e^{-a^2} \sum_{k=0}^{n-2} \frac{(-1)^k}{k!} \right). \]  

(95)

The first negative moment of \( \lambda \) can also be evaluated through (93), given as

\[ E\left[ \frac{1}{\lambda} \right] = 2 \left( \sqrt{2\epsilon} \right)^{n-1} \Gamma (n-1) \left( 1 + e^{-a^2} \sum_{k=0}^{n-2} \frac{(-1)^k}{k!} \right). \]  

(97)

Substituting (95) and (97) in Proposition 3, we obtain the desired result.

\section*{APPENDIX E}

\section*{PROOF OF LEMMA 4}

We note that \( E[I] \) can be evaluated straightforwardly since \( \hat{g}_k \sim CN (0, \beta_k) \). Let \( \Omega = \sum_{k=1,k \neq j}^{N} |\hat{g}_k|^2 \) for convenience. To obtain \( E[I^j] \), it is necessary to establish the pdf of \( \Omega \). Note that \( \hat{g}_k \sim CN (0, \beta_k) \), and the sum of \( |\hat{g}_k|^2 \) follows a generalized chi-squared distribution, given as [43, Theorem 4]

\[ f_\Omega (x; N-1, \beta_1, \ldots, \beta_j-1, \beta_j+1, \ldots, \beta_N) = \frac{e^{-\frac{x}{\beta_j}}}{\beta_j^\frac{N-1}{2} \Gamma \left( \frac{N-1}{2} \right)} \prod_{k=1,k \neq j}^{N} \beta_k \left( \frac{1 - \frac{\beta_k}{\beta_j}}{\sqrt{\beta_j}} \right). \]  

(99)

The first negative moment of \( \Omega \) can be derived as (98), where \( (a) \) is derived by using [44, Eq. (3.352.4)].
\[ E[\lambda] = \frac{e^{-\sum_{i,j} e_{ij}}}{m^{\prod_{k<l} (\epsilon_i - \epsilon_k)}} \sum_{j=1}^{m} \sum_{i=1}^{m} \frac{m_n^{-2} 2^{2n-m-n-2}}{K+1} \sum_{j=1}^{m} e^{x_j} Q_{n-m+2k+1,n-m} \left( \sqrt{2e_{ij}}, 0 \right) U_{i,j} \]

\[ E[I^{-1}] = \sum_{k=1}^{N} \sum_{k \neq k'} \rho_{max}^\beta_k \prod_{l=1,l \neq k,l \neq j}^{N} \left( 1 - \frac{\beta_{l'}}{\rho_{max}} \right) \int_{0}^{\infty} \frac{1}{1 + x} e^{-\frac{\pi}{m} x} \, dx \]

\[ e^{\frac{-\pi}{2m}} \]

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